



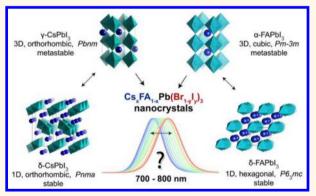
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# **Exploration of Near-Infrared-Emissive Colloidal** Multinary Lead Halide Perovskite Nanocrystals Using an Automated Microfluidic Platform

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Supporting Information

ABSTRACT: Hybrid organic-inorganic and fully inorganic lead halide perovskite nanocrystals (NCs) have recently emerged as versatile solution-processable light-emitting and light-harvesting optoelectronic materials. A particularly difficult challenge lies in warranting the practical utility of such semiconductor NCs in the red and infrared spectral regions. In this context, all three archetypal A-site monocationic perovskites—CH<sub>3</sub>NH<sub>3</sub>PbI<sub>3</sub>, CH(NH<sub>2</sub>)<sub>2</sub>PbI<sub>3</sub>, and CsPbI<sub>3</sub>—suffer from either chemical or thermodynamic instabilities in their bulk form. A promising approach toward the mitigation of these challenges lies in the formation of multinary compositions (mixed cation and mixed anion). In the case of multinary colloidal NCs, such as quinary Cs, FA1-, Pb-(Br<sub>1-v</sub>I<sub>v</sub>)<sub>3</sub> NCs, the outcome of the synthesis is defined by a



complex interplay between the bulk thermodynamics of the solid solutions, crystal surface energies, energetics, dynamics of capping ligands, and the multiple effects of the reagents in solution. Accordingly, the rational synthesis of such NCs is a formidable challenge. Herein, we show that droplet-based microfluidics can successfully tackle this problem and synthesize Cs<sub>x</sub>FA<sub>1-x</sub>PbI<sub>3</sub> and Cs<sub>x</sub>FA<sub>1-x</sub>Pb(Br<sub>1-x</sub>I<sub>y</sub>)<sub>3</sub> NCs in both a time- and cost-efficient manner. Rapid in situ photoluminescence and absorption measurements allow for thorough parametric screening, thereby permitting precise optical engineering of these NCs. In this showcase study, we fine-tune the photoluminescence maxima of such multinary NCs between 700 and 800 nm, minimize their emission line widths (to below 40 nm), and maximize their photoluminescence quantum efficiencies (up to 89%) and phase/chemical stabilities. Detailed structural analysis revealed that the Cs<sub>x</sub>FA<sub>1-x</sub>Pb(Br<sub>1-y</sub>I<sub>y</sub>)<sub>3</sub> NCs adopt a cubic perovskite structure of FAPbI, with iodide anions partially substituted by bromide ions. Most importantly, we demonstrate the excellent transference of reaction parameters from microfluidics to a conventional flaskbased environment, thereby enabling up-scaling and further implementation in optoelectronic devices. As an example,  $Cs_rFA_{1-r}Pb(Br_{1-r}I_r)_3$  NCs with an emission maximum at 735 nm were integrated into light-emitting diodes, exhibiting a high external quantum efficiency of 5.9% and a very narrow electroluminescence spectral bandwidth of 27 nm.

KEYWORDS: perovskites, microfluidics, nanocrystals, formamidinium, quantum dots, halides

ead halide perovskites (LHP) of the APbX<sub>3</sub> type, where A can be methylammonium (MA, CH<sub>3</sub>NH<sub>3</sub><sup>+</sup>), formamidinium (FA,  $CH_3(NH_2)_2^+$ ), inorganic cations (Cs<sup>+</sup>, Rb<sup>+</sup>), or a

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mixture thereof and X is a halide (Br, I, and their mixtures), have attracted enormous attention after they were recognized as efficient thin-film absorber materials for photovoltaics, with power conversion efficiencies of up to 22.7%. 1-4 Recently, there has been a surge of studies on the nanoscale counterparts of these perovskites and, in particular, on colloidal nanocrystals (NCs), 5-12 which hold great promise as versatile photonic sources for displays, lighting, light-emitting diodes (LEDs), and lasers<sup>13–16</sup> and as light harvesters for solar cells and photodetectors. 1,17 Unlike other forms of APbX<sub>3</sub> perovskites—bulk single and polycrystals, thin films, and other substrate-grown structures—colloidal NCs have a particular set of advantages, foremost of which is their versatile solution processability and miscibility with other materials, as well as access to quantum-size effects and their facile surface- and shape-engineering. <sup>18</sup> To date, essentially all work on colloidal APbX3 NCs has concentrated on those compositions, which emit in the visible region of the electromagnetic spectrum (between 400 and 700 nm), with CsPbBr3 and FAPbBr3 NCs, exhibiting green photoluminescence (PL, between 500 and 550 nm), being by far the most popular targets. This can be attributed to the high durability of these bromides, as compared to other compositions in the APbX<sub>3</sub>

In the current study, we focus our attention on multinary perovskite NCs of the quinary composition  $Cs_rFA_{1-r}Pb(Br_{1-r}I_r)_3$ , with a view to controlling their optical properties in the near-infrared (700–800 nm). It should be noted that there is extensive literature regarding red and near-infrared emissive colloidal MAPbI<sub>3</sub> NCs, refs 7 and 19-23 being representative examples. However, MAPbI<sub>3</sub> NCs suffer from severe instabilities caused by humidity, light, and heat and eventually decompose into CH<sub>3</sub>NH<sub>2</sub>, PbI<sub>2</sub>, HI, and  $I_2^{22,24,25}$  Fully inorganic CsPbX<sub>3</sub> NCs (X = Cl, Br, and I), synthesized by a hot injection method, exhibit tunable emission between 410 and 700 nm and high PL quantum efficiencies (50-90%). Presently, CsPbX<sub>3</sub> NCs are the focus of attention with respect to their chemical engineering (i.e., identification of precursors, growth kinetics, shape-control, postsynthetic reactivity, and up-scaling), <sup>20,26–45</sup> surface chemistry, <sup>8,46–50</sup> photophysics (single-dot spectroscopy, lasing, *etc.*), <sup>51–59</sup> and applications in television displays, <sup>6,60–62</sup> light-emitting devices, <sup>50,63–68</sup> and solar cells.<sup>17</sup> The usability of red-emissive CsPbX<sub>3</sub> NCs is, however, strongly limited by the phase instability of the 3D polymorphs of CsPbI<sub>3</sub>. For similar reasons, iodide-rich CsPb(Br/I)<sub>3</sub> compositions are also unstable. As for FA-based analogues, which benefit from the higher chemical stability of the FA ion (compared to that of MA<sup>+</sup>), one-pot colloidal syntheses have led to the development of stable and highly emissive FA-doped CsPbI<sub>3</sub> NCs (PL peak at ca. 690 nm; ca. 10% FA) and FAPbI<sub>3</sub> NCs (PL peak at ca. 780 nm), 69 with cubic or nearly cubic NC shapes and mean particle sizes between 10 and 20 nm. Interestingly, unlike colloidal FAPbI<sub>3</sub> NCs, bulk FAPbI<sub>3</sub> is completely unstable due to a phase transition from semiconductive perovskite into a yellow, nonperovskite phase (an observation that will be further discussed in the Results and Discussion section of this study). A formidable challenge with iodide-based NCs lies in discovering the compositional space and related synthesis parameters that would provide for continuous coverage of the 700-800 nm spectral range with narrow PL line widths and without compromising the chemical durability of the NCs.

Our choice of the quinary composition,  $Cs_xFA_{1-x}Pb(Br_{1-y}I_y)_3$ , was motivated by previous thin-film and bulk single-crystal studies, <sup>70–74</sup> which indicate that the concomitant incorporation of Cs and Br into the FAPbI<sub>3</sub> lattice stabilizes the 3D perovskite

phase and allows for the compositional tuning of band-gap energy through the formation of mixed-halide solid solutions. With monohalide FAPbI $_3$  and FA $_{0.1}$ Cs $_{0.9}$ PbI $_3$  NCs, the only way to access the entire 700–800 nm range is to exploit quantum-size effects and reduce NC size to the 3–8 nm range. Such small NCs are highly labile morphologically, and their PL characteristics are consequently broad and unstable. <sup>76</sup>

In the case of NCs, accessing suitable reaction parameters for the formation of stable Cs<sub>x</sub>FA<sub>1-x</sub>PbX<sub>3</sub> NCs with tunable emission maxima in the range of 700–780 nm is simply not possible with flask reactions. Specifically, in a recent study on the formation of mixed-anion FAPbX<sub>3</sub> NCs, it was reported that variation in the Br–I content led solely to two stable compositions, one emitting at 680 nm and the other at 760 nm. This is due to the fact that in NCs the ability to form certain compositions is governed not only by the thermodynamics of the mixedion phases but also by the NC surfaces (*i.e.*, surface energy and ligand binding) and chemical equilibria with the precursors in the solution. These factors greatly expand relevant parametric space beyond the mixing ratios of the Cs, FA, Pb, and halide precursors, thus making thorough exploratory synthesis and optimization virtually impossible with conventional flask-based techniques.

Recently, we showcased the potential of droplet-based microfluidics combined with online absorption and PL spectroscopy in discovering the optimal synthesis parameters for ternary and quaternary (mixed-anion) CsPbX<sub>3</sub> NCs<sup>77</sup> and FAPbX<sub>3</sub><sup>78,79</sup> NCs. The next logical step, which is pursued in this work, is the exploration of NCs of higher compositional complexity, i.e.,  $Cs_xFA_{1-x}Pb(Br_{1-y}I_y)_3$  NCs. In such cases, the use of highthroughput automated microfluidic reactors becomes an absolute necessity for rapid and detailed experimentation. Using microfluidics, we herein demonstrate the formation of  $Cs_rFA_{1-r}Pb(Br_{1-r}I_r)_3$  NCs with PL peak tuning between 690 and 780 nm and narrow PL line widths (expressed as full width at half-maximum, fwhm). Characterization by X-ray diffraction (XRD) and elemental analysis pointed to the incorporated quantities of Cs (0.1–1.2% with respect to FA) and Br (10–18% with respect to I). The obtained  $Cs_xFA_{1-x}Pb(Br_{1-y}I_y)_3$  NCs were approximately cubic in shape, with a mean size of 8-13 nm. The reaction parameters were fully transferrable from our microfluidic platform (pL-nL scale) to conventional reaction flasks (mL scale). This allows us to demonstrate the optoelectronic utility of these NCs by fabricating near-infrared LEDs with an external quantum efficiency (EQE) up to 5.9% at 735 nm.

# **RESULTS AND DISCUSSION**

Formability of APbl<sub>3</sub>. To obtain red to near-infrared emissive  $Cs_xFA_{1-x}Pb(Br_{1-y}I_y)_3$  compositions, the structural chemistry of these perovskites must be considered: only those polymorphs of APbI<sub>3</sub> compounds that retain three-dimensional (3D) corner-shared interconnections of lead halide octahedra are potent semiconductors. The extended electronic structure in lead halide perovskites arises from the periodic 3D network of PbX<sub>6</sub> octahedra. Low-dimensionality polymorphs of the same composition exhibit higher (by at least 1 eV) and often indirect band gaps, typically on the order of 300-500 nm. 3D polymorphs of all archetypal ternary APbI3 compounds have stability issues. MAPbI3 is chemically unstable, whereas the 3D polymorphs of CsPbI<sub>3</sub> and FAPbI<sub>3</sub> are thermodynamically metastable and undergo transitions into 1D polymorphs (Figure 1). 76,80-85 Interestingly, the thin-film and nanoscale forms of CsPbI<sub>3</sub> and FAPbI<sub>3</sub> exhibit extended but still finite stability in their 3D

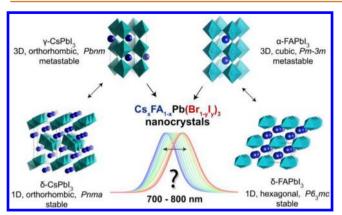


Figure 1. Formabilities of the 3D and 1D polymorphs of CsPbI<sub>3</sub> and FAPbI3 compounds and the goal of this study: near-infrared emissive LHP NCs. The PbI<sub>6</sub> octahedra of α-FAPbI<sub>3</sub> NCs are assembled in a 3D cubic metastable lattice, which spontaneously converts into a 1D hexagonal version (nonluminescent) at room temperature. In the case of CsPbI3, the PbI6 octahedra of FAPbX3 NCs are assembled in a 3D orthorhombic metastable lattice ( $\gamma$ -phase), which eventually converts at room temperature into a 1D orthorhombic  $\delta$ -phase (nonluminescent). The goal of this study is highlighted with a question: can high-throughput microfluidic screening identify the existence of stable multinary  $Cs_xFA_{1-x}Pb(Br_{1-y}I_y)_3$  phases in the form of colloidal NCs, which cover the PL region of 700-800 nm, i.e., in-between ternary 3D phases (CsPbI3 and FAPbI3)? We note that bulk α-FAPbI<sub>3</sub> emits at 840 nm and γ-CsPbI<sub>3</sub> emits at 710 nm, whereas their NC counterparts are commonly reported to emit at  $\leq$ 700 and  $\leq$ 780 nm, respectively.<sup>6,8-12,69</sup> It is also noted that the space groups reported for the  $\gamma$ - and  $\delta$ -phases of CsPbI<sub>3</sub> do not differ (while their structures manifestly do), as they can easily be interconverted by simple axis permutations. We used the original Phnm and Pnma for the  $\gamma$ - and  $\delta$ -forms, respectively, to maintain consistency with past literature.

polymorph forms (from days to several months), primarily due to surface effects.  $^{6,32,69,86-88}$ 

Thermodynamic instability is caused by the Cs and FA ions being, respectively, slightly too small and too large for the voids in between  $PbI_6$  octahedra. This has been broadly discussed in the literature in terms of the Goldsmith tolerance factor (t) and octahedral factor  $(\mu)$ , which describe the optimal dense packing of charged ions in an ideal cubic 3D perovskite. Page Mixing larger and smaller ions at the A-site is a powerful strategy for adjusting the geometric fitness of the A-cation in the void space of the 3D lead halide framework, thereby improving its phase stability. A high entropy of mixing is also considered to be a stabilizing factor. In fact, the best phase stabilities and optoelectronic performances have been reported for thin films with such mixed A-site occupations: FA/MA,  $^{94-97}$  Cs/MA,  $^{98}$  Cs/FA,  $^{70-73}$  Cs/MA/FA, or even Rb/Cs/MA/FA.

Toward Multinary APbX<sub>3</sub> Perovskites: The Goal of This Investigation. Covering the desired spectral range of 700–800 nm requires broad mapping in Cs-FA perovskites, because 3D CsPbI<sub>3</sub> and FAPbI<sub>3</sub> exhibit band gaps of 1.75 eV (710 nm) and 1.48 eV (840 nm), respectively. From the viewpoint of structural chemistry, it remains unclear which crystal structures will be adopted or are stable when Cs:FA and Br:I ratios are adjusted simultaneously, as the parent Cs and FA compounds have different 3D polymorph structures (Figure 1). We recently found that for bulk single crystals, in accordance with previous studies on thin films starting with α-FAPbI<sub>3</sub>, one can concomitantly introduce up to 15% Cs and 30% Br, while maintaining the same cubic crystal structure as α-FAPbI<sub>3</sub>.

As for colloidal NCs, we recently reported the synthesis of a mixed-cationic composition,  $Cs_{0.9}FA_{0.1}PbI_3$ , with the same crystal structure as orthorhombic 3D  $\gamma$ -CsPbI<sub>3</sub>, via a flask-based process. Multiple variations of the Cs:FA reagent ratio and synthesis temperatures led to the same composition (with *ca*. 10% FA). From these studies, it can be surmised that the exploration of complex compositions is prohibitively difficult *via* flask-based syntheses. This is because with one iteration of only one parameter per synthesis (each lasting several hours when conducted manually) several years will be required to properly map compositional space and other factors, such as the effects of ligands, solvents, and solvation equilibria.

Experimental Design and Combinatorial Strategy. Microfluidic reactors allow for the addition of multiple reagents in a user-defined manner, rapid thermal and mass transfer, and quantitative kinetic investigation of reactions, thus defining an ideal medium for preparing semiconductor NCs with well-defined morphologies and physicochemical properties. 100–103 In addition, the advances in robust microfluidic configurations, real-time detection methods, 104–111 continuous purification 112,113 and ligand-exchange 114 systems, and optimization algorithms 115–117 make microfluidic reactors ideal for the detailed investigation of rapid and complex reaction kinetics 77,78,106,118 and for the discovery of multicomponent semiconductor NCs. 77,78,119

Herein, we modified and applied a previously developed microfluidic platform that had been used for the synthesis and real-time characterization of binary chalcogenide NCs<sup>106</sup> and CsPbX<sub>3</sub> NCs.<sup>77</sup> This platform (see Figure 2 and associated description in the Methods section) incorporates a multiphase microfluidic reactor with integrated PL and absorption detection to rapidly screen reaction conditions. The controlled injection of precursor solutions (Cs-oleate, FA-oleate, PbX<sub>2</sub>, and PbY<sub>2</sub>) and carrier fluid is performed in an automated manner (using syringe pumps), allowing for the formation of nanoliter droplets using a seven-port manifold and efficient mixing of precursors (in  $\sim 300 \text{ ms}$ ).<sup>77</sup> In the case of  $Cs_xFA_{1-x}PbX_3$  (X = Br and I) NCs, we defined four interdependent molar ratios, which were adjusted during synthesis: FA/Pb, Cs/Pb, Br/I, and Cs/FA. In this report, the latter is presented as the Cs percentage relative to the FA content. A tube-based microfluidic reactor allows for rapid heating of the droplets (within a few hundred ms), along with the real-time extraction of PL and absorption characteristics at various reaction times (0.1-20 s) and temperatures (25-130 °C). Additional characterization of the synthesized NCs by transmission electron microscopy (TEM) and XRD was conducted in an off-line manner, by collecting a sufficient quantity of the sample during synthesis at a fixed set of reaction conditions (see the Methods section for details). The optimized parameters were then transferred to conventional flask-based reactions.

**Synthesis of Cs<sub>x</sub>FA<sub>1-x</sub>PbI<sub>3</sub> NCs.** In our previous study, we showed that FAPbI<sub>3</sub> NCs can grow even at room temperature, <sup>78</sup> while CsPbI<sub>3</sub> NCs with a 3D phase are formed at temperatures above 100 °C. <sup>77</sup> Accordingly, we decided that analyzing the temperature range suitable for the formation of  $Cs_xFA_{1-x}PbI_3$  NCs was an important initial task. Figure 3a–c report the variation in emission line width and PL peak at 25–130 °C (FA/Pb = 9.3, Cs/Pb = 0.3, and %Cs = 3.0). The formation of FAPbI<sub>3</sub> NCs takes place at room temperature, which can be inferred by the emergence of the PL peak at 792 nm, consistent with a previous flask-based synthetic study, in which 15 nm FAPbI<sub>3</sub> NCs exhibited a PL peak at 780 nm. <sup>69</sup> Higher temperatures led to a rapid increase in the band-gap energy, which we ascribe to the

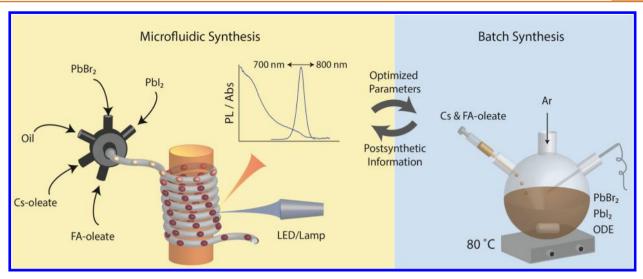


Figure 2. (Left) Illustration of the segmented-flow reaction platform equipped with online PL and absorbance modules for the synthesis and real-time monitoring of  $Cs_xFA_{1-x}PbX_3$  perovskite NCs. The microfluidic platform allows for a systematic and independent variation of precursor molar ratios, such as Cs/Pb, FA/Pb, Cs/FA, and Br/I, growth times (determined by the flow rate and tube lengths), and temperature. Droplets are generated by adjusting the flow rates of the carrier phase ( $50-200~\mu L/min$ ) and that of the dispersed phase ( $1.2-50~\mu L/min$ ). (Right) Illustration of a typical flask-based hot-injection synthesis of  $Cs_xFA_{1-x}PbX_3$  NCs. Overall, synthesis optimization was performed by mutual information exchange between flask-based experimentation (identification of suitable precursors, solvents, and capping ligands) and microfluidics (optimization of the reaction parameters). The optimized reaction parameters were successfully transferred from microfluidics back into flask reactors, followed by up-scaling and additional postsynthetic characterization (XRD, electron microscopy, and stability tests).

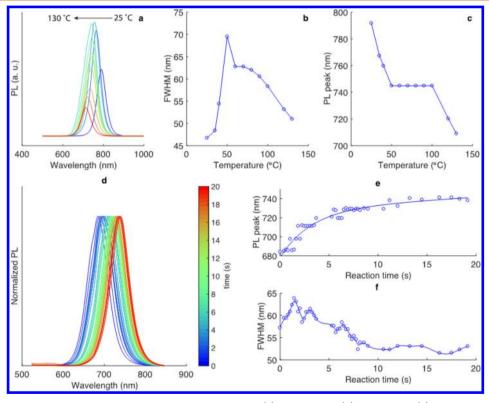


Figure 3. Microfluidic synthesis of  $Cs_xFA_{1-x}PbI_3$  NCs. Variation in the (a) PL spectra, (b) fwhm, and (c) PL maximum as a function of temperature for  $Cs_{0.03}FA_{0.97}PbI_3$  NCs (with the variation in the Cs/FA molar ratio indicated). Other parameters were as follows: FA/Pb = 9.3, Cs/Pb = 0.3, and reaction time = 10 s. (d-f) Temporal evolution of the normalized online PL spectra, PL maxima, and fwhm of  $Cs_{0.02}FA_{0.98}PbI_3$  NCs at 80 °C.

incorporation of Cs. The possibility of smaller NC sizes (quantum dots 3–10 nm in diameter) causing larger band gaps can be discounted using postsynthesis TEM images, with all NCs obtained in this study for growth times greater than 7 s being 15–20 nm in size. Size evolution occurs very quickly and over the

course of several seconds. From 50 to 90  $^{\circ}$ C, the PL peak remained stable at 740–745 nm (Figure 3c), but with a gradually decreasing fwhm. Higher synthesis temperatures (>110  $^{\circ}$ C) resulted in PL peaks closer to 700 nm, most likely due to the formation of ternary CsPbI<sub>3</sub> NCs with or without minimal

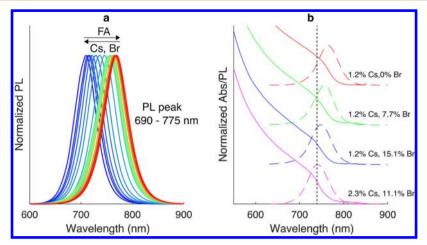


Figure 4. (a) PL spectra of colloidal  $Cs_xFA_{1-x}Pb(Br_{1-y}I_y)_3$  NCs synthesized using the microfluidic platform and exhibiting composition-tunable band-gap energies between 690 and 780 nm with fwhm values of 40–65 nm and (b) representative online PL and in online absorption spectra at different quantities of  $Cs^+$  and  $Br^-$  in the reaction mixture.

incorporation of FA ions. Accordingly, we concluded that the  $50-90~^{\circ}\text{C}$  range is ideal for compositional engineering purposes.

The high speed of formation of  $Cs_{0.02}FA_{0.98}PbI_3$  NCs is on par with FAPbX<sub>3</sub> (see ref 78 and Figure S1) and CsPbI<sub>3</sub> (reported previously),<sup>77</sup> taking only a few seconds to stabilize the PL maximum at 740 nm and the fwhm at 52 nm (Figure 3d–f). Offline optical characterization after synthesis indicates that there is no subsequent growth or other form of evolution in the  $Cs_xFA_{1-x}PbI_3$  NCs (see Figure S2). Based on these observations, further rapid automated screening, at a rate of 100 adjustments per synthesis parameter per hour, was carried out with reaction times of at least (and typically) 7 s.

The efficiency of Cs incorporation is expected to depend not only on temperature but also on the solvation conditions and Cs/FA ratio and to some extent on the Cs/Pb and FA/Pb ratios. Because these relationships are not fully and rationally predictable when the equilibrium constants and involved energies (lattice energies for all compositions, solvation energies, surface energies, and ligand binding energies) are not completely known, they were tested in this study in a combinatorial fashion (Figures S3–S6). In brief, for excess FA (by a factor of 6 with respect to Pb), a narrow fwhm can be obtained (Figure S3). At much higher FA/Pb ratios (>13), the crystal phase of the NCs tends to change from black to yellow within hours of synthesis. The TEM images revealed severe morphological irregularities in such NCs, which were in the form of large populations of micrometer-sized needles and rods (Figure S4). Furthermore, the PL tunability of Cs<sub>x</sub>FA<sub>1-x</sub>PbI<sub>3</sub> NCs was limited to Cs/Pb ratios lower than 2 and Cs loadings of  $\leq 10\%$  (Figure S3). Outside this window, the PL maxima were always in the range of 680-700 nm, suggesting the formation of pure CsPbI<sub>3</sub> NCs or their mixtures with other compositions (seen as multi-Gaussian PL lines; see Figure S3e, for Cs content equal to 23.3%) or perhaps even mixtures of various shapes. The TEM images (Figure S4, S5) illustrate how the three interlinked ratios affect the morphology of the synthesized Cs<sub>x</sub>FA<sub>1-x</sub>PbI<sub>3</sub> NCs. Progressive addition of 0.3-5.2% Cs+ (with respect to the FA content) continuously tunes the PL maximum from 758 to 710 nm (Figure S6), while maintaining a narrow fwhm in the range of 48-55 nm. Although the XRD results suggest that mixed perovskites adopt a structure similar to that of pure FAPbI<sub>3</sub> (Figure S7), the majority of Cs<sub>x</sub>FA<sub>1-x</sub>PbI<sub>3</sub> NC samples exhibited

low colloidal and chemical stabilities, except those samples with both low FA/Pb ratios ( $\leq$ 7) and low Cs loadings (up to 2%).

Synthesis of  $Cs_xFA_{1-x}Pb(Br_{1-y}I_y)_3$  NCs. The addition of Br into the system was then explored as a way of stabilizing crystal structure and tuning the PL maximum in the desired range (700-800 nm). The operating temperature and reaction times were similar to those used for the synthesis of  $Cs_xFA_{1-x}PbI_3$  NCs. As the addition of a second halide further increases the complexity of the synthetic system and can shift the product equilibrium toward other perovskite compositions, parametric screening must be performed with caution. We therefore broadly explored the influence of the interdependent molar ratios of Cs/Pb, FA/Pb, and Br content on the optical properties and stability of the Br/I mixtures (as illustrated in Figures S8–S10). In brief, such a combinatorial study revealed that the parametric zones of the FA/Pb and Cs/Pb molar ratios, able to tune the PL peak between 690 and 780 nm (while maintaining a satisfactory fwhm), were 2.5-6.0 and 0.01-0.04, respectively (Figure S9). In addition, Br loading of up to 15% leads to a linear blue shift of the emission band at all FA/Pb molar ratios (Figure S10). A key message here is that in nearly all optimized compositions Br addition does not alter the emission line width or emission intensity, suggesting that the synthesized NCs have stable optical characteristics. Furthermore, an increase in Br loading over 25% will deliver perovskite NCs with emission energies in the range of 650-720 nm. However, such an increase in Br loading can trigger the formation of other perovskite structures, such as  $FAPb(Br_{1-\nu}I_{\nu})_3$  NCs, due to excess FA-oleate in the reaction system (Figure S11).

Figure 4a presents selected PL spectra in the range of 690 to 775 nm, with fwhm in the range of 45–65 nm. Such precise PL tuning is achieved through a systematic variation of all three interlinked molar ratios within their refined parametric zones. In particular, variation in Cs (in the range of 0.2–5.2%) and Br (between 0% and 15% of the total halide concentration) content leads to a blue shift in the in-line (*i.e.*, postheating and when the reaction was quenched) absorption and PL spectra (Figure 4b). Most importantly, the incorporation of up to 15% Br into the structure of  $Cs_xFA_{1-x}PbI_3$  NCs increased the period of stability of the  $Cs_xFA_{1-x}PbI_3$  NCs from several hours to several weeks (Figure S12).

Transfer to Flask-Based Synthesis. To assess whether the optimal parameters can be transferred to conventional

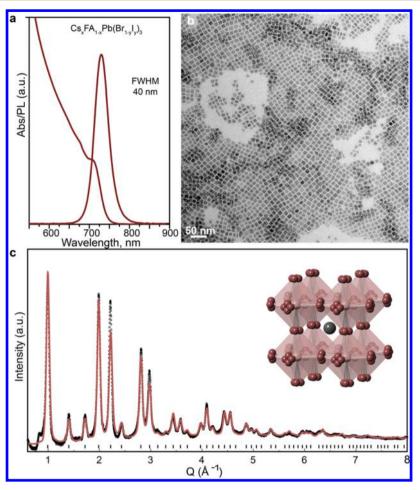


Figure 5. Optical absorption and PL spectra of  $Cs_xFA_{1-x}Pb(Br_{1-y}I_y)_3$  NCs synthesized in conventional flask reactors, exhibiting a fwhm of 40 nm. (b) Bright-field scanning TEM (STEM) image of  $Cs_xFA_{1-x}Pb(Br_{1-y}I_y)_3$  NCs. (c) Synchrotron XRD pattern (black) and best fit (red,  $2\theta$  range of 0.5–130°;  $\lambda$  = 0.563 729 Å) for  $Cs_xFA_{1-x}Pb(Br_{1-y}I_y)_3$  NCs, yielding a refined lattice parameter (a = 6.3296 Å) and the anionic composition. The inset illustrates the cubic perovskite structure of  $Cs_xFA_{1-x}Pb(Br_{1-y}I_y)_3$  NCs (space group  $Pm\overline{3}m$ , with y = 0.87 and x = 0), in which the perovskite framework consists of  $PbX_6$  units sharing the octahedral corners; the  $X^-$  anions are disordered in four equivalent positions.

flask reactors, we carried out hot-injection synthesis of both  $Cs_xFA_{1-x}PbI_3$  and  $Cs_xFA_{1-x}Pb(Br_{1-y}I_y)_3$  NCs (further details are provided in the Methods section). Briefly, to synthesize NCs with the nominal composition of  $Cs_{0.01}FA_{0.99}Pb(Br_{0.11}I_{0.89})_3$ (Figure 5), PbI<sub>2</sub> (55 mg, 0.12 mmol) and PbBr<sub>2</sub> (5 mg, 0.014 mmol) were suspended in 1-octadecene (ODE, 4.6 mL), heated to 60 °C, and then dried under vacuum for 30 min. Subsequently, the reaction mixture was heated to 110 °C in a nitrogen environment, followed by the addition of dried solvents: oleylamine (OLA, 0.5 mL) and oleic acid (OA, 1.0 mL). Once the PbI<sub>2</sub> dissolved, the reaction mixture was cooled to 80 °C. At this point, a mixture of FA oleate (4.8 mL) and Cs oleate (1.2 mL) stock solutions was injected into the reaction flask. After 5 s, the reaction was quenched. The crude solution was then centrifuged and the supernatant was discarded. The precipitate was dissolved in hexane, and the resulting solution was centrifuged once again, after which the supernatant and precipitate were separated. The particles from both fractions, supernatant and precipitate, were further washed to remove excess organic ligands (see further details in the Methods section). This synthesis procedure yielded nearly cubic NCs with a PL peak at 730 nm (Figure 5a-c), a PL fwhm of 40 nm (after isolation and purification), and a PL quantum yield (QY) of 80-89%. QY drops to ca. 50% in the solid-state form (NC film). We note that the larger fwhm values detected in-line in microfluidics can be

attributed to size-fractioning that occurs during isolation and purification.

Crystal Structure. To uncover the structural details of  $Cs_xFA_{1-x}Pb(Br_{1-y}I_y)_3$  NCs, synchrotron X-ray total scattering measurements were performed (using an octane solution of NCs in a quartz capillary, Figure 5c) at the X04SA-MS4 Powder Diffraction Beamline of the Swiss Light Source (Paul Scherrer Institute, Villigen, CH). 120 A combined Rietveld and total scattering approach based on the Debye scattering equation (DSE, accounting for structure, size, and anisotropic morphology)<sup>121</sup> was used for structural and microstructural characterization of  $Cs_xFA_{1-x}Pb(Br_{1-y}I_y)_3$  NCs (details are provided in the Supporting Information). The analysis results suggest a cubic structure in which X<sup>-</sup> anions are disordered in four equivalent sites around the equilibrium position (inset of Figure 5c), similar to the disorder seen in FAPbI<sub>3</sub> and FAPbBr<sub>3</sub> NCs. 69,122 The atomic displacement parameters (in the form of the Debye-Waller factor) were refined for all atoms; the anomalously high values for halides in the unsplit cubic arrangement suggest local structural disorder. The graphical outcomes of the DSE-based analysis of Cs<sub>x</sub>FA<sub>1-x</sub>Pb(Br<sub>1-y</sub>I<sub>y</sub>)<sub>3</sub> NCs are summarized in Figures 5c, S13, and S14. In a similar manner to CsPbX<sub>3</sub> NCs, <sup>12</sup> the peak positions slightly deviate from the cubic metric. Nevertheless, the hybrid NCs investigated here exhibited a different kind of structural defectiveness, which needs further investigation.

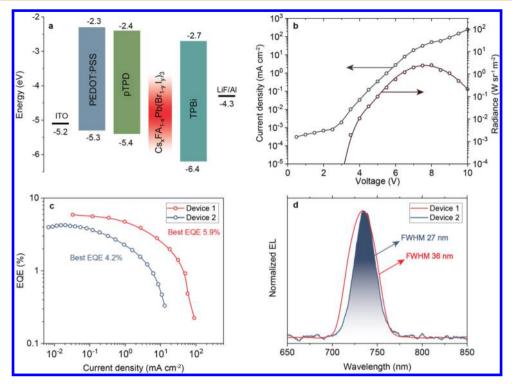


Figure 6. (a) Energy diagram of LED devices with  $Cs_xFA_{1-x}Pb(Br_{1-y}I_y)_3$  NCs as emissive layers. (b) Current density and radiance *versus* voltage characteristics of device 1. (c) External quantum efficiency *versus* current density characteristics shown for devices 1 and 2. (d) Narrowest EL spectra of device 1 and device 2.

In addition, to further validate our findings, we compared the experimental data with a Pnma orthorhombic structure model (Figure S15). 123 The Pnma structure does not reproduce the experimental peak intensities; at the same time, new peaks appear in a simulation, without having a counterpart in the experimental data, thus supporting the analysis presented in Figure 5c. To investigate the substitutional disorder between the FA<sub>1-x</sub>/Cs<sub>x</sub> and  $Br_{1-v}/I_{\nu}$  couples, the corresponding site occupancy factors (s.o.f.) were refined by the conventional Rietveld method, and the resulting values (x = 0; y = 0.87) were kept fixed during DSEbased modeling. The Br/I substitution value is consistent with the PL peak position and with that estimated by X-ray fluorescence (XRF). The Cs quantity was too low to be detected by X-ray techniques; however, a small quantity of this cation (<5%) might be present in the crystal structure. Concerning the structurally refined s.o.f. value of Br (0.13), an even more robust determination (s.o.f = 0.10) was independently derived by adopting Vegard's law, which correlates the refined lattice parameter and anionic composition of the mixed halide, Cs,FA1-xPb- $(Br_{1-\nu}I_{\nu})_3$  (for  $x \approx 0$  and a = 6.3296 Å), with the two end members of the same series (FAPbI<sub>3</sub>, a = 6.3639 Å, and FAPbBr<sub>3</sub>, a = 6.0042 Å),  $^{69,122}$  as depicted in Figure S14. Additionally, using a standard benchtop energy dispersive XRF instrument and a calibration mixture of Pb(NO<sub>3</sub>)<sub>2</sub> and KBr in a 5:1 ratio, the Br fraction in the title compound was calculated to be 0.16(3), corroborating the presented estimates. In addition, to assess whether the quinary compositions undergo phase separation into respective ternary compounds, we report the pattern of a mixture of FAPbBr<sub>3</sub> (13% w/w) and FAPbI<sub>3</sub> (87% w/w) cubic phases and compare it to a solid solution model (Figure S16). The result indicates the absence of phase segregation. Overall, we highlight that FA<sup>+</sup> cations, fully or nearly fully occupying the A-site of the 3D perovskite framework, systematically favor the formation of a cubic structure [FAPbI<sub>3</sub>, FAPbBr<sub>3</sub>, and the herein studied

 $Cs_xFA_{1-x}Pb(Br_{1-y}I_y)_3$ ], whereas in  $CsPbX_3$  NCs, the smaller Cs ions favor the tilting of octahedra and the formation of orthorhombic  $\gamma$ -phases, even after incorporation of up to 10% FA.<sup>69</sup>

Thermal stability of quinary NCs and, for comparison, FAPbI<sub>3</sub> NCs was evaluated using thermogravimetry and differential scanning calorimetry (DSC); see Figure S17. The decomposition of both kinds of NCs occurs in a few steps. Both show a most pronounced DSC feature at 335 °C that can be attributed to the decomposition of FA cations, indicating that both materials are of similar thermal stability.

**Light-Emitting Diodes.**  $Cs_rFA_{1-r}Pb(Br_{1-r}I_r)_3$  NCs were used to construct LEDs with the device structure depicted in Figure 6a. These LEDs were fabricated by spin coating poly(3,4ethylenedioxythiophene)-poly(styrenesulfonate) (PEDOT:PSS) on prepatterned indium tin oxide (ITO)-covered glass substrates, followed by the spin coating of  $poly(N_iN'-bis(4-butylphenyl)-$ *N,N'*-bisphenylbenzidine (poly-TPD) and an NC emissive layer. Later, 50 nm of 2,2,2"-(1,3,5-benzinetriyl)tris(1-phenyl-1Hbenzimidazole) (TPBi), 1.2 nm of lithium fluoride (LiF), and 100 nm of aluminum (Al) were sequentially evaporated on top of the emissive layer. Two different devices (1 and 2) were tested, with device 2 containing TOPAS polymer in the NC layer (TOPAS = cyclic olefin copolymer from TOPAS Advanced Polymers GmbH). LED performance was characterized by measuring current density and radiance as a function of the voltage applied between the ITO anode and Al cathode (Figure 6b) and by measuring the electroluminescence (EL) spectrum (Figure 6d). The current density of device 1 steadily increased from 8 ×  $10^{-4}$  mA cm<sup>-2</sup> at 2 V to over 100 mA cm<sup>-2</sup> at 10 V. The radiance surpassed 10<sup>-4</sup> W sr<sup>-1</sup> m<sup>-2</sup> at just above 3 V and increased up to 3.9 W sr<sup>-1</sup> m<sup>-2</sup> at 7.5 V. The radiance, current density, and EL spectrum were then used to calculate the EQE, which describes the number of out-coupled photons per number of injected electrons. The EQE dependence on current density is shown in

Figure 6c, with a device 1 peak EQE of 5.9% at 0.1 mA cm<sup>-2</sup>  $(0.2\% \text{ at } 100 \text{ mA cm}^{-2})$  and a device 2 peak EQE of 4.2%. The peak EQE and turn-on voltage of >3 V are in line with our previous work, in which we investigated FAPbI<sub>3</sub> LEDs emitting at 772 nm. 69 However, in the case of LEDs incorporating  $Cs_xFA_{1-x}Pb(Br_{1-y}I_y)_3$  NCs, the EQE stays above 2.5% for currents up to 10 mA cm<sup>-2</sup>. The peak radiance is also higher with quinary NCs (3.9 versus 1.54 W sr<sup>-1</sup> m<sup>-2</sup>).<sup>69</sup> While the low efficiency roll-off below 10 mA cm<sup>-2</sup> indicates a good electron hole balance in the emission layer 124 and excellent carrier/ exciton confinement,  $^{125}$  the turn-on voltage of over 3 V,  $\sim 1.3$  V higher than expected from the band-gap energy of the  $Cs_xFA_{1-x}Pb(Br_{1-\nu}I_{\nu})_3$  NCs (~1.7 eV), suggests that further optimization of the device architecture is needed in order to reduce the charge-injection barrier(s). The EL spectra (Figure 6d) revealed a sharp peak near 735 nm with a narrowest fwhm of 37 nm at 9 V for device 1 and as low as 27 nm at 8.5 V for device 2. Figures S18 and S19 present the evolution of EL spectra with increasing the voltage, for both devices. Even though device 1 showed the best EQE value of 5.9% and peak radiance of  $3.9\,\mathrm{W\,sr^{-1}\,m^{-2}}$ , device 2 (best EQE of 4.2%) showed a distinctly different behavior: EL peak narrowing at higher voltages (Figure S19), from an initial fwhm of 40 nm to 27 nm at 8.5 V. In device 2, the fwhm of the EL spectrum is strikingly narrower than that in the NC PL spectra in the solution (fwhm = 40 nm) and in films (fwhm = 52 nm); see Figure S20 for comparison. Together with the red shift of the EL peak position, this might indicate an efficient energy transfer between the NCs. 126 As a plausible scenario, the applied voltage might induce anion migration and hence alter the energy band gaps of NCs within the layer. 127 The emission might then occur through the channeling of excitation into the specific population of NCs. To the best of our knowledge, the EL fwhm of 27 nm (62 meV) is the narrowest among those reported for the red (and near-IR) perovskite LEDs. Figure S21 shows typical transient properties of the NC LED at a constant voltage of 4.5 V. Significant growth of the current density from 0.1 mA cm<sup>-2</sup> to over 5 mA cm<sup>-2</sup> within seconds is observed. This might again support the possibility of the ionic rearrangement within the NC film that improves charge injection, similar to the light-emitting electrochemical devices. The overall LED device lifetime (to reach half of the maximum EL intensity) at 4.5 V bias is about half a minute. For comparison, other perovskite LEDs in this wavelength range are those utilizing methylammonium lead iodide (MAPbI<sub>3</sub>) thin films treated with *n*-butylammonium iodide (EQE of 10.4%),  $^{129}$  quasi-2D perovskites (EQE of 8.8%), <sup>126</sup> and multiple quantum wells (EQE of 11.7%).<sup>130</sup>

# **CONCLUSIONS**

Herein, we have described the combinatorial synthesis of highly luminescent and stable  $Cs_xFA_{1-x}Pb(Br_{1-y}I_y)_3$  NCs with emission and absorption spectra between 690 and 780 nm. Using a microfluidic platform, we identified the compositional parametric zones for fine-tuning optical properties, while retaining emission line widths in the range of 45–65 nm (before isolation, size selection, and purification). Importantly, microfluidic parameters can be readily transferred to conventional flask-based synthetic processes used by the perovskite research community. The PL fwhm can be further refined to ~40 nm as a result of the size selection occurring during the isolation and purification of NCs. Characterization by synchrotron X-ray scattering indicates a cubic structure for the  $Cs_xFA_{1-x}Pb-(Br_{1-x}I_y)_3$  NCs, wherein solid mixed-anion solutions can be

clearly confirmed. The distribution of Cs ions remains rather controversial, yet its addition to the reaction mixture is of paramount importance in improving the long-term stability of NCs. The successful fabrication of NC-based LED devices with EQEs as high as 5.9% indicates the excellent chemical durability of  $Cs_xFA_{1-x}Pb(Br_{1-y}I_y)_3$  NCs. Future studies should focus on understanding the interesting observations of the slow transient response of the LEDs and narrow EL spectral line width of 27 at 735 nm.

The synthesis of multinary lead halide perovskite (LHP) NCs might open opportunities for the broad deployment of these nanomaterials in photovoltaics and other optoelectronic devices. For instance, these infrared-active NCs are highly desirable for applications in single-junction or tandem photovoltaics, where NC colloids can be employed as inks for depositing active absorbing layers. <sup>1,17</sup> In this regard, in contrast to conventional molecular LHP solutions used as inks, the ease of compositional engineering and facile ligand removal exhibited by the currently developed quinary NCs, followed by low-temperature sintering for recrystallization into continuous thin films, suggests numerous possibilities. Alternatively, or rather orthogonally, methods of surface coating for maintaining the quantum-size effects may enable thin-film or quantum-dot-like photovoltaic devices <sup>17,131</sup> and photodetectors. <sup>132</sup>

#### **METHODS**

**Materials.** Cesium carbonate ( $Cs_2CO_3$ , Aldrich, 99.9%), formamidinium acetate (Sigma-Aldrich, 99.9%), lead bromide (PbBr<sub>2</sub>, ABCR, 98%), lead iodide (PbI<sub>2</sub>, ABCR, 99.999%), 1-octadecene (90%), oleic acid (Sigma-Aldrich, 90%), and oleylamine (Acros, ≥ 96%) were used as the reagents. Galden PFPE fluid was purchased from Blaser Swisslube AG. Patterned indium tin oxide (ITO)-covered glass substrates were purchased from Lumtech. The hole injection material PEDOT:PSS was purchased from Heraeus (CLEVIOS VP AI 4083), while the hole transport material poly-TPD was procured from Lumtech and the electron transport material TPBi was supplied by e-Ray Optoelectronics. The electron injection material LiF was purchased from Acros Organics, and Al pellets were purchased from Kurt J. Lesker Co. Ltd. TOPAS polymer (cyclic olefin copolymer) was received from TOPAS Advanced Polymers GmbH. All the materials for LED production were used as received without any further purification.

Microfluidic Synthesis. Various concentrations of precursor solutions (see Supporting Information for details) were used depending on the experimental purpose (sample collection, absorption measurements, PL measurements). Precision syringe pumps (neMESYS, Cetoni GmbH, Germany) were used to inject the dispersed phase (PbX<sub>2</sub>, FA-oleate, and Cs-oleate precursor solutions) and the carrier fluid (Galden fluorinated fluid, Blaser Swisslube AG, Germany) toward a manifold (Manifold Assay 7 Port 10-32 Std, Upchurch Scientific, Germany) to form a segmented flow of droplets. The injection manifold and the syringes carrying the precursor solutions were connected through polytetrafluoroethylene tubing (i.d. 250  $\mu$ m, o.d. 1/16 in., Upchurch Scientific, Germany) using polyether ether ketone fingertight fittings (F-127, Upchurch Scientific, Germany). The carrier fluid was transferred to the manifold via fluorinated ethylene propylene tubing (i.d. 750  $\mu$ m, o.d. 1/16 in., Upchurch Scientific, Germany). Typical flow rates were between 80 and 100  $\mu$ L min<sup>-1</sup> for the carrier phase and between 0.1 and 50  $\mu$ L min<sup>-1</sup> for the precursors. The chemical payload of the formed droplets can be tuned in a precise and rapid fashion by continuously varying the precursor volumetric flow rates. The formed droplets containing the reaction mixture were subsequently directed through perfluoroether tubing (i.d. 500  $\mu$ m, o.d. 1/16 in., Upchurch Scientific, Germany) coiled around a copper heating rod (diameter = 1.5 cm) to allow initiation of the NC-forming reaction and online detection of the formed perovskite NCs. The overall reaction time was kept constant in all experiments by ensuring a

constant tubing length between the point where the tubing enters the heating rod and the detection volume.

Online Photoluminescence Measurements. A 375 nm LED (M375L3-Mounted LED, Thorlabs, Germany) was used as an excitation source for PL measurements. The collimated beam was directed toward a dichroic beam splitter (Multiphoton LP-Strahlenteiler HC 375 LP, AHF, Germany) and then focused into the microfluidic channel using an aspheric lens (A240TM, f=8.0 mm, NA 0.50, Thorlabs, Germany). Emission originating from the microfluidic channel was collected by the same lens, passed through the dichroic beam splitter, and coupled *via* a 10× objective (RMS10X, NA 0.25, Thorlabs, Germany) to a fiber spectrometer (QE 65000, Ocean Optics, UK) *via* a 2 m long multimode fiber with a core diameter of 400  $\mu$ m (QP400-2-UV—vis, Ocean Optics, UK). The spectrometer incorporated a 20  $\mu$ m entrance slit, a 600 lines/mm grating, and a 2048-pixel detector. The spectrometer was operated between 350 and 1100 nm, and data were recorded using an integration time between 50 and 100 ms.

Online Absorbance Measurements. Absorbance measurements were conducted after the heating stage, where the reaction mixture flows through a high-purity perfluoralkoxy capillary (1/16 in. o.d.,  $500 \mu m$  i.d., IDEX Health & Science, USA). The in-line absorbance spectrometer consists of a fiber-coupled halogen lamp (HL-2000 HP, Ocean Optics, UK) and a fiber-coupled spectrometer (AvaSpec ULS2048 Starline, Avantes, USA). The spectrometer was operated between 200 and 1100 nm, and data were recorded using an integration time of 100 ms.

Flask Synthesis: Preparation of Formamidinium Oleate Stock Solution. Formamidinium acetate (3.765 mmol, 0.392 g, Aldrich, 99%) was loaded into a 50 mL three-neck flask along with ODE (18 mL) and OA (12 mL). The reaction mixture was degassed three times at room temperature, heated to  $100\,^{\circ}\mathrm{C}$  in a nitrogen environment, maintained at that temperature until the reaction is complete, and then cooled to room temperature. The resulting solution was stored in a glovebox.

Flask Synthesis: Preparation of a Cesium Oleate Stock Solution. Cesium carbonate (0.015 mmol, 5 mg) was loaded into a 25 mL three-neck flask along with ODE (10 mL) and OA (0.625 mL). The reaction mixture was degassed three times at room temperature, heated to 120 °C in a nitrogen atmosphere, maintained at that temperature until the reaction was complete, and cooled to room temperature. The resulting solution was stored in a glovebox.

Flask Synthesis of  $Cs_{0.01}FA_{0.99}Pb(Br_{0.11}I_{0.89})_3$  NCs. In a 25 mL three-necked flask, PbI<sub>2</sub> (55 mg, 0.12 mmol, Sigma-Aldrich) and PbBr<sub>2</sub> (5 mg, 0.014 mmol, Sigma-Aldrich) were suspended in ODE (4.6 mL), heated to 60 °C, and then dried under vacuum for 30 min. Subsequently, the reaction mixture was heated to 110 °C in a nitrogen atmosphere, followed by the addition of dried solvents: OLA (0.5 mL, Strem) and OA (1.0 mL, Aldrich). Once PbI<sub>2</sub> was dissolved, the reaction mixture was cooled to 80 °C. At this point, a mixture of the formamidinium oleate (4.8 mL) and cesium oleate (1.2 mL) stock solutions was injected into the reaction flask. After another 15 s, the reaction mixture was cooled using a water-ice bath. The crude solution was centrifuged at 12 100 rpm for 7 min, and the supernatant discarded. The precipitate was dissolved in hexane (250  $\mu$ L), and the resulting solution was centrifuged again at 10 000 rpm for 3 min. The supernatant and precipitate were separated, and 150 µL of hexane was added to the supernatant. This fraction was labeled as "SN". The precipitate was dissolved in toluene (1.0 mL) and centrifuged at 3500 rpm for 2 min to get rid of large NCs, with the resultant sample being labeled as "P". Particles from both fractions were washed again to remove excess organic ligands.

Washing of  $Cs_{0.01}FA_{0.99}Pb(Br_{0.11}I_{0.89})_3$  SN NCs. To 100  $\mu$ L of a hexane solution of NCs were added hexane (100  $\mu$ L), toluene (200  $\mu$ L), and methyl acetate (530  $\mu$ L). The solution was centrifuged at 13 400 rpm for 3 min and redissolved in hexane or toluene (PL at ~730 nm).

Washing of  $Cs_{0.01}FA_{0.99}Pb(Br_{0.11}I_{0.89})_3$  P NCs. Methyl acetate (0.65 mL) was added to a toluene solution of fraction "P" and centrifuged for 3 min at 10 000 rpm. The obtained precipitate was dissolved in toluene, hexane, or octane (PL at ~750 nm).

**Offline Characterization.** Ultraviolet—visible (UV—vis) absorbance spectra were recorded using a Jasco V770 spectrometer in transmission

mode. Photoluminescence spectra were recorded using a Fluoromax iHR 320 Horiba Jobin Yvon spectrofluorimeter equipped with a PMT detector. The excitation wavelength was 400 nm, and the excitation source was a 450 W xenon lamp. The measured intensities were corrected to account for the spectral response of the detector. Powder XRD patterns were recorded using a powder diffractometer (STOE STADI P) with Cu K $\alpha$ 1 radiation. The diffractometer was operated in transmission mode with a germanium monochromator and a silicon strip detector (Dectris Mythen). TEM images were captured using a JEOL JEM-2200FS microscope operated at 200 kV. Quantitative XRF measurements were conducted using a benchtop Minipal 2 PANalytical spectrometer with polycarbonate films supporting dry colloids or powders and a Cr X-ray tube operating at a maximum power of 30 W. Thermal analysis (thermogravimetry and differential scanning calorimetry) was performed using a Netzsch Simultaneous thermal analyzer (STA 449 F5 Jupiter). A powdered sample (6–10 mg) was placed in an alumina crucible and heated under Ar gas flow (50 mL/min) to 800 °C (10 °C/min). NC solutions in hexane were predried in small alumina beakers at room temperature.

**Quantum Yield Measurements.** To measure the relative PL quantum yield of the NC solution, dilute solutions of NCs in toluene and dye standards (Rhodamine 6G in ethanol and zinc phthalocyanine in benzene) were prepared in 10 mm optical path length cuvettes, ensuring an absorbance of approximately 0.1 at either 488 or 633 nm. PL quantum yields were calculated according to

$$\Phi_f^i = \frac{F^i f_s n_i^2}{F^s f_i n_s^2} \Phi_f^s$$

where  $\Phi^i_f$  and  $\Phi^s_f$  are the PL QYs of the sample and standard, respectively;  $F^i$  and  $F^s$  are the integrated areas of the sample and standard spectra, respectively;  $f_i$  and  $f_s$  are the absorption factors of the sample and standard ( $f = 1 - 10^{-1}$ , where A is the absorbance), respectively;  $n_i$  and  $n_s$  are the refractive indices of the sample and standard, respectively. <sup>133</sup>

Absolute quantum yield measurements of the films and NC solutions were performed using a Hamamatsu Quantaurus QY spectrometer (C11347-11) equipped with an integrating sphere. The excitation peak wavelength was 450 nm.

Fabrication of LED Devices. Initially, ITO substrates were rinsed with a mixture of deionized water and detergent solution. Subsequently, substrates were sonicated for 20 min in acetone and isopropyl alcohol. To enhance wettability, substrates were treated with an oxygen plasma for 10 min. An aqueous solution of PEDOT:PSS was spin-coated at 4000 rpm for 30 s, after which the ITO substrates were annealed on a hot plate under ambient conditions for 30 min at 130 °C. Subsequently, they were transferred into a nitrogen-filled glovebox for the deposition of subsequent layers. Poly-TPD was spin coated at 1000 rpm from a 2 mg/mL chlorobenzene solution and annealed for 20 min at 120 °C. For device 1 a colloidal suspension of  $Cs_xFA_{1-x}Pb(Br_{1-y}I_y)_3$  NCs in hexane (7 mg/mL) was spin coated at 2000 rpm; for device 2 the NC solution was mixed with TOPAS polymer (0.5 mg/mL) and spin coated at 2000 rpm. Subsequently, the substrates were transferred into a vacuum chamber at  $10^{-7}$  mbar, where 50 nm of TPBi (electron transport layer), 1.2 nm of LiF, and 100 nm of Al (cathode) were evaporated through a shadow mask at evaporation rates of 0.5, 0.1, and 2 A/s, respectively. An active pixel area of 16 mm<sup>2</sup> was determined by the overlap of ITO and Al. All devices were measured under ambient conditions without encapsulation.

**LED Performance Characterization.** The J-V-L characteristics of the fabricated LEDs were measured under ambient conditions using a Keysight 2902b source measurement unit and a calibrated photodiode (FDS1010-CAL, Thorlabs). The size of the photodiode  $(10 \times 10 \text{ mm}^2)$  is much larger than that of the active pixel size  $(4 \times 4 \text{ mm}^2)$  of the LEDs. The EQEs of the fabricated LEDs were calculated from the known EL spectra of the LEDs and photodiode sensitivity, <sup>134</sup> while the radiance was calculated assuming a Lambertian emission profile. EL spectra were recorded using a CCS200 CCD spectrometer (Thorlabs) and a PR-655 (Photoresearch) spectroradiometer. LED transient optoelectronic properties were measured with a Keysight 2902b source measurement unit and amplified photodiode (PDA36A-EC, Thorlabs) with 20  $\mu$ s resolution.

Synchrotron X-ray Total Scattering Measurements. The experimental synchrotron X-ray total scattering data of colloidal  $Cs_xFA_{1-x}(Br_{1-y}I_y)_3$  NCs  $(FA = [HC(NH_2)_2]^+)$  NCs in octane were collected at the X04SA-MS4 Powder Diffraction Beamline of the Swiss Light Source (Paul Scherrer Institute, Villigen, CH) using a certified quartz capillary (0.5 mm in diameter). The operational beam energy was set to 22 keV ( $\lambda$  = 0.563 729 Å) and accurately determined using a silicon powder standard (NIST 640d,  $a_0 = 0.543 123(8)$  nm at 22.5 °C). Data were collected in the 0.5–130°  $2\theta$  range using a single-photoncounting silicon microstrip detector (MYTHEN II). Total scattering patterns with air background, empty glass capillary, and pure solvent were independently collected under the same experimental conditions and properly subtracted from the sample signal. Transmission coefficients of the sample- and solvent-loaded capillaries were also measured and used for angle-dependent absorption correction. Inelastic Compton scattering was added as an additional model component during data analysis. For DSE-based modeling, an angular range of 3-120° was used.

Temperature-dependent measurements were also performed in the range of 98-348 K using a temperature-controlled  $N_2$  stream fluxing over the capillary. No phase transitions were observed in the explored temperature range.

#### **ASSOCIATED CONTENT**

# S Supporting Information

The Supporting Information is available free of charge on the ACS Publications website at DOI: 10.1021/acsnano.8b01122.

Additional schematic illustration, preparation of precursors, synthesis procedure in a flask, TEM images, XRD patterns, and the PL and absorption spectra of  $Cs_xFA_{1-x}PbX_3$  NCs (PDF)

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#### **Notes**

The authors declare no competing financial interest.

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